## Hunting glueballs with BESIII

Complutense University of Madrid, April 3rd, 2019
A. Rodas


JPAC/BESIII: A workshop on Theory-Experiment collaboration.

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## This work: Motivation

- In this talk: Ongoing phenomenological analysis on spectroscopy
- Ordinary hadrons $\rightarrow$ Boring!!

- Not so ordinary $\rightarrow$ Not today!



## This work: Motivation

- In this talk: Glueball spectroscopy?
- Hybrid $\rightarrow$ part of this talk



## . Glueball

## This work: Motivation

- Glueball expected at around $1.5-2 \mathrm{GeV}$.
- $J^{P C}=0^{++} \rightarrow$ lightest glueball candidate(s).

| $J^{P C}$ | Other $J$ | $r_{0} m_{G}$ | $m_{G}(\mathrm{MeV})$ |  |
| :--- | :---: | :---: | :---: | :---: |
| $0^{++}$ |  | $4.21(11)(4)$ | $1730(50)(80)$ |  |
| $2^{++}$ |  | $5.85(2)(6)$ | $2400(25)(120)$ |  |
| $0^{-+}$ |  | $6.33(7)(6)$ | $2590(40)(130)$ |  |
| $0^{*++}$ |  | $6.50(44)(7)^{\dagger}$ | $2670(180)(130)$ |  |
| $1^{+-}$ |  | $7.18(4)(7)$ | $2940(30)(140)$ |  |
| $2^{-+}$ |  | $7.55(3)(8)$ | $3100(30)(150)$ |  |
| $3^{+-}$ |  | $8.66(4)(9)$ | $3550(40)(170)$ |  |
| $0^{*-+}$ |  | $8.88(11)(9)$ | $3640(60)(180)$ |  |
| $3^{++}$ | $6,7,9, \ldots$ | $8.99(4)(9)$ | $3690(40)(180)$ |  |
| $1^{--}$ | $3,5,7, \ldots$ | $9.40(6)(9)$ | $3850(50)(190)$ |  |
| $2^{*-+}$ | $4,5,8, \ldots$ | $9.50(4)(9)^{\dagger}$ | $3890(40)(190)$ |  |
| $2^{--}$ | $3,5,7, \ldots$ | $9.59(4)(10)$ | $3930(40)(190)$ |  |
| $3^{--}$ | $6,7,9, \ldots$ | $10.06(21)(10)$ | $4130(90)(200)$ |  |
| $2^{+-}$ | $5,7,11, \ldots$ | $10.10(7)(10)$ | $4140(50)(200)$ |  |
| $0^{+-}$ | $4,6,8, \ldots$ | $11.57(12)(12)$ | $4740(70)(230)$ |  |



PDG status

- Glueball expected at around $1.5-2 \mathrm{GeV}$.
- Three different candidates measured close by.

| $f_{0}(1370)\left[{ }^{[]]}\right.$ | $I^{G}\left(J^{P C}\right)=0^{+}\left(0^{+}+\right)$ |  |
| :---: | :---: | :---: |
| Mass $m=1200$ to 1500 MeV <br> Full width $\Gamma=200$ to 500 MeV |  |  |
| $\mathrm{f}_{0}(1370)$ DECAY MODES | Fraction ( $\left.\Gamma_{i} / \mathrm{r}\right)$ | $p(\mathrm{MeV} / \mathrm{c})$ |
| $\pi \pi$ | seen | 672 |
| $4 \pi$ | seen | 617 |
| $4 \pi^{0}$ | seen | 617 |
| $2 \pi^{+} 2 \pi^{-}$ | seen | 612 |
| $\pi^{+} \pi^{-2} \pi^{0}$ | seen | 615 |
| $\rho \rho$ | dominant | $\dagger$ |
| $2(\pi \pi)_{S \text {-wave }}$ | seen | - |
| $\pi(1300) \pi$ | seen | $\dagger$ |
| $a_{1}(1260) \pi$ | seen | 35 |
| $\eta \eta$ | seen | 411 |
| $K \bar{K}$ | seen | 475 |
| $K \bar{K} n \pi$ | not seen | $\dagger$ |
| $6 \pi$ | not seen | 508 |
| $\omega \omega$ | not seen |  |
| $\gamma \gamma$ | seen | 685 |
| $e^{+} e^{-}$ | not seen | 685 |

$$
f_{0}(1500)^{[n]} \quad \quad I^{G}\left(J^{P C}\right)=0^{+}\left(0^{++}\right)
$$

Mass $m=1504 \pm 6 \mathrm{MeV} \quad(\mathrm{S}=1.3)$
Full width $\Gamma=109 \pm 7 \mathrm{MeV}$

| $\boldsymbol{f}_{\mathbf{0}} \mathbf{( 1 5 0 0 )}$ DECAY MODES | Fraction $\left(\Gamma_{i} / \Gamma\right)$ | Scale factor | $\rho$ <br> $(\mathrm{MeV} / \mathrm{c})$ |
| :--- | :--- | ---: | ---: |
| $\pi \pi$ | $(34.9 \pm 2.3) \%$ | 1.2 | 740 |
| $\pi^{+} \pi^{-}$ | seen |  | 739 |
| $2 \pi^{0}$ | seen |  | 740 |
| $4 \pi$ | $(49.5 \pm 3.3) \%$ | 1.2 | 691 |
| $4 \pi^{0}$ | seen |  | 691 |
| $2 \pi^{+} 2 \pi^{-}$ | seen |  | 686 |
| $2(\pi \pi)_{S \text {-wave }}$ | seen |  | - |
| $\rho \rho$ | seen |  | $\dagger$ |
| $\pi(1300) \pi$ | seen |  | 143 |
| $a_{1}(1260) \pi$ | seen |  | 217 |
| $\eta \eta$ | $(5.1 \pm 0.9) \%$ | 1.4 | 515 |
| $\eta \eta^{\prime}(958)$ | $(1.9 \pm 0.8) \%$ | 1.7 | $\dagger$ |
| $K \bar{K}$ | $(8.6 \pm 1.0) \%$ | 1.1 | 568 |
| $\gamma \gamma$ | not seen |  | 752 |

${ }_{1}{ }^{G}\left(J^{P C}\right)=0^{+}\left(0^{++}\right)$
Mass $m=1723_{-5}^{+6} \mathrm{MeV} \quad(\mathrm{S}=1.6)$
Full width $\Gamma=139 \pm 8 \mathrm{MeV} \quad(\mathrm{S}=1.1)$
$f_{0}(1710)$ DECAY MODES

| $K \bar{K}$ | seen | $p(\mathrm{MeV} / c)$ |
| :--- | :--- | ---: |
| $\eta \eta$ | seen | 706 |
| $\pi \pi$ | seen | 665 |
| $\omega \omega$ | seen | 851 |

## Consensus?

- The glueball is expected to be predominant in either the $f_{0}(1500)$ or the $f_{0}(1710)$.
- Not much of a consensus $\rightarrow$ V. Mathieu et al. Int.J.Mod.Phys. E18 (2009) 1-49.
- Recent years $\rightarrow$ not much of an improvement.
- $f_{0}(1500) \rightarrow 0.89|g g\rangle$ Giacosa et al. Phys.Rev. D72 (2005) 094006.
- $f_{0}(1710) \rightarrow 0.93|g g\rangle$ Albaladejo-Oller Phys.Rev.Lett. 101 (2008) 252002.


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## Data: Glueball "rich" experiments

- Pomeron collisions
- $p \bar{p}$ anihilation

- $J / \psi$ radiative decays considered the golden channel for glueballs.



## Data: BESIII $J / \psi \rightarrow \gamma \pi \pi$

- Data on $J / \psi \rightarrow \gamma \pi \pi$ half a million events.


- 3 prominent $f_{0}$ 's with similar couplings.
- The $2^{++} E 1$ partial wave is dominated by the $f_{2}(1270)$.


## Data: BESIII $J / \psi \rightarrow \gamma K \bar{K}$

- Another 3 prominent $f_{0}$ 's


- The couplings are fairly different, with a way more prominent $f_{0}(1710)$.
- The $2^{++} E 1$ partial wave is dominated by the $f_{2}^{\prime}(1525)$.


## Data: BESIII $J / \psi$

- How many $f_{0}$ do we have here?



- Is the coupling of the $f_{0}(1710)$ greater $\rightarrow$ glueball hint?


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## Method

- Based on ar et al. Phys.Rev.Lett. (2019), A.Jackura et al. Phys.Lett.B (2018)
- Peripheral production $\Rightarrow$ factorization of the pomeron $\Rightarrow$ $\operatorname{Ima}(s)=\rho(s) t^{*}(s) a(s)$.
- Amplitude built around
$t(s)=\frac{N(s)}{D(s)}$ method
$\Rightarrow a(s)=p^{2} q \frac{n(s)}{D(s)}$.
- They are smooth polynomials $n(s)=\sum_{j} a_{j} w^{j}(s)$, where $w(s)=\frac{s}{s+s_{0}}$.



## Method

- $N(s)$ and $n(s)$ are process dependent, they have only left hand cuts.
- $\mathrm{D}(\mathrm{s})$ has a right hand cut, altogether $t(s)$ has the correct analytic structure.

- By adding this discontinuity over the RHC one could go to the direct continuous Riemann sheet.



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## Coupled channel

- $\eta^{\left({ }^{\prime}\right)} \pi$ coupled channel up to 2 GeV .
- We use a K-matrix approach with a Chew-Mandelstam phase space.

$$
\begin{aligned}
D^{J}(s)_{k i} & =\left(K^{J}(s)^{-1}\right)_{k i}-\frac{s}{\pi} \int_{s_{k}}^{\infty} d s^{\prime} \frac{\rho\left(s^{\prime}\right) N_{k i}^{J}\left(s^{\prime}\right)}{s^{\prime}\left(s^{\prime}-s-i \varepsilon\right)} \\
\rho N_{k i}^{J}\left(s^{\prime}\right) & =\delta_{k i} \frac{\lambda^{J+1 / 2}\left(s^{\prime}, m_{\eta^{\prime},}^{2}, m_{\pi}^{2}\right)}{\left(s^{\prime}+s_{L}\right)^{2 J+1+\alpha}} \\
K_{k i}^{J}(s) & =\sum_{R} \frac{g_{k}^{J, R} g_{i}^{J, R}}{m_{R}^{2}-s}+c_{k i}^{J}+d_{k i}^{J} s
\end{aligned}
$$

- Just 1 K-matrix pole for the P-wave.


## Coupled channel analysis

- We use an average of 6 parameters for each figure.
- $\chi^{2} \approx 1.3$, no significant deviation for any partial wave.
- 1 K-matrix pole produces 2 different peaks for the P -wave $\rightarrow$ 300 MeV distance.




## Poles

- Statistical uncertainties calculated through bootstraping
- $m\left(a_{2}\right)=1306.0 \pm 0.8 \pm 1.3 \mathrm{MeV}$ $\Gamma\left(a_{2}\right)=114.4 \pm 1.6 \pm 0.0$ MeV
- $m\left(a_{2}^{\prime}\right)=1722 \pm 15 \pm 67 \mathrm{MeV}$ $\Gamma\left(a_{2}^{\prime}\right)=247 \pm 17 \pm 63 \mathrm{MeV}$
- $m\left(\pi_{1}\right)=1564 \pm 24 \pm 86 \mathrm{MeV}$ $\Gamma\left(\pi_{1}\right)=492 \pm 54 \pm 102 \mathrm{MeV}$.
- All systematics (diferent LHC masses, numerator models ...) included.

$J / \psi \rightarrow \gamma m_{1} m_{2}$
- Slightly different kinematics

- Left hand cut $\rightarrow s=0 \mathrm{GeV}$.
- Ima $(s)=\rho(s) t(s)^{*} a(s)$
- $t(s) \rightarrow \pi \pi, K \bar{K}$ scattering.
$\pi / K$


$\pi / K$


## Coupled-channel scenario

- Fit from 1 GeV to $2.5 \mathrm{GeV}, \chi^{2} \approx 1.5$.
- Interested in the $f_{0}$.
- Coupled channel between just $\pi \pi$ and $K \bar{K}$.
$\mathrm{J} / \psi \rightarrow \gamma \pi \pi$ (S wave)




## Complex plane

- We use the analytical properties of the parameterization $\rightarrow$ complex plane continuation.

- $m\left(f_{0}(1500)\right)=1460 \mathrm{MeV}$
- $m\left(f_{0}(1710)\right)=1800 \mathrm{MeV}$
- $m\left(f_{0}(210)\right)=1970 \mathrm{MeV}$
$\Gamma\left(f_{0}(1500)\right)=85 \mathrm{MeV}$.
$\Gamma\left(f_{0}(1710)\right)=190 \mathrm{MeV}$.
$\Gamma\left(f_{0}(1710)\right)=490 \mathrm{MeV}$.


## Scalar poles

- Complex plane plots


- Few "spurious" poles, all far from real axis


## Improvements: $\sigma$ description

- Tree level ChPT $\rightarrow T^{0}(s, t) \propto \frac{s-M_{\pi}^{2} / 2}{f_{\pi}^{2}}$.
- PCAC $\rightarrow$ Adler zero $\rightarrow K_{k i}^{J} \propto\left(s-s_{A}\right)$.

$$
K_{k i}^{J}(s)=\frac{s-s_{A}}{s}\left[\sum_{R} \frac{g_{k}^{J, R} g_{i}^{J, R}}{m_{R}^{2}-s}+c_{k i}^{J}+d_{k i}^{J}\right],
$$

- Dispersive Adler zero located at $s_{A}=85 \mathrm{MeV}^{2}$.GKPRY Phys.Rev.D (2012)



## Improvements: $\sigma$ description

- Wrong behavior at low energies.
- Even the Adler zero is not sufficent to directly accommodate the $\sigma$ pole.
- Bump produced by the background+phase space
- Solution $\rightarrow$ including $\pi \pi$ data.




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## Future Prospects: Dispersive $\pi \pi$ description

- Can be done using $T=V+V G T$, Ropertz-Kubis-Hanhart Eur.Phys.J. c78 (2018) $\bar{B}_{s}^{0} \rightarrow J / \psi \pi \pi / K \bar{K}$.
- Omnés-like factorization of the

$$
\sigma \rightarrow T=T_{\sigma}+\Omega\left[1-V_{R} \Sigma\right]^{-1} \Omega^{t} .
$$

- $\operatorname{Im} \Omega_{i j}=\left(T_{0}\right)_{i m}^{*} \sigma_{m} \Omega_{m j}$.
- Used to accommodate the low energy $\pi \pi$ dispersive input.




## Summary

- Implementation of a coupled-channel formalism $\rightarrow$ In progress.
- Description of the features of $3 f_{0}$ 's.
- Description of the $\sigma \rightarrow \ln$ progress.
- Inclusion of the dispersive $\pi \pi$ result $\rightarrow$ Next step.


## Thank you for your attention!

